

Improved Interface Properties and Dielectric Breakdown in Recessed AlGaN/GaN MOS-HEMTs Using HfSiO_x as Gate Dielectric

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Abstract—In this letter, we report optimized transport properties in gate recessed enhancement-mode GaN MOS-HEMTs by incorporating silicon into atomic layer deposited gate dielectric HfO₂. Compared with commonly used HfO₂ gate dielectric, the interface trap density can be reduced by nearly an order of magnitude and the fixed oxide traps inside are reduced to almost half using the high-quality passivation of HfSiO_x. The MOS-HEMTs based on HfSiO_x exhibit a threshold voltage of 1.5 V, excellent subthreshold swing of 65 mV/dec, and a high on–off ratio of 3×10^{10} . The incorporation of silicon in HfO₂ can also increase the dielectric breakdown property with maximum gate electric field of 2.85 MV/cm for a 10-year time-dependent gate dielectric breakdown lifetime, which is 36% higher than pure HfO₂. The maximum breakdown voltage of HfSiO_x MOS-HEMT is 742 V, which is 30% higher than HfO₂ MOS-HEMT.

Index Terms—Atomic layer deposition, dielectric breakdown, GaN MOS-HEMT, recess gate, interface trap density.

I. INTRODUCTION

NORMALLY-OFF GaN HEMTs are essential for power conversion to maintain system safety and reliability. However, the high density 2DEG at the heterointerfaces keeps the channel normally-on without applying an external electrical field. Typical methods to achieve normally-off HEMTs include fluoride plasma injection, [1], [2], p-type gate [3], [4], partially or fully recessed gate barrier [5]–[7], and a mixture of these mentioned methods [8], [9]. Although threshold voltage can be increased, many of these gate-related modifications will introduce side effects such as high gate leakage and low gate swing. To simplify the gate driver complexity and avoid device failure from gate voltage overshoot, high quality gate dielectrics are necessary not only to suppress gate leakage, but also to maintain 2DEG’s inherent high mobility, especially

in the case of recess gate structure. Indeed, 2DEG under the recessed area typically suffers from scatterings from poor dielectric/GaN interface, leading to deteriorated gate reliability which poses limits to the commercialization of recessed gate MOS-HEMTs [10], [11].

Numerous dielectrics have been used in recessed gate HEMTs along with various deposition technologies, such as Al₂O₃ [5], [12], MgCaO [13], LaHfO_x [14], SiON [15], SiN_x [6] and multi-layer dielectrics like SiN_x/HfO₂ [16], SiO₂/LaAlO₃ [17] and AlN/Al₂O₃ [18]. High- κ gate dielectric HfO₂ stands out among many dielectrics with its high dielectric constant compared with other widely used dielectrics such as SiN_x, SiO₂, Al₂O₃, and AlN. With its higher dielectric constant, HfO₂ MOS-HEMTs can achieve much more efficient electrostatic control. However, HfO₂ on GaN MOS-HEMTs suffers from high leakage current because of insufficient barrier height [19], which will deteriorate device performance via gate leakage. Recent studies have shown the incorporation of silicon into Al₂O₃ can enhance breakdown strength and improve interface properties [20], [21]. In order to improve the dielectric performance of HfO₂, we investigated the inclusion of silicon into HfO₂.

In this letter, we report recessed gate E-mode GaN MOS-HEMTs using silicon doped HfO₂ dielectric and study the effects on device performance from traps between dielectric and channel interface and traps inside dielectrics. HfSiO_x can reduce interface trap density and the fixed oxide traps inside the dielectric, which maintains 2DEG’s high mobility property under the recessed gate while achieving normally-off operation. Meanwhile, the incorporation of silicon in HfO₂ can boost the breakdown properties of the dielectric, and thus the breakdown voltage of the device.

II. DEVICE FABRICATION

Fig. 1(a) shows the cross-section schematic view of gate recessed GaN MOS-HEMTs. The epitaxial structure consists of a 4 μ m buffer layer, a 1 nm AlN insertion layer and a 20 nm Al_{0.25}Ga_{0.75}N barrier layer grown on (111) silicon substrate. The fabrication process started with mesa isolation. Gate recess was performed using low-power Cl₂/BCl₃ plasma etching with an etching depth \sim 20 nm and a smooth etching angle, as shown in Fig. 1(b). Atomic force measurement was used to evaluate the uniformity of the recess surface with a small root mean square roughness of around 0.464 nm. The Ohmic contact regions were deposited with Ti/Al/Ni/Au metal stack, followed by rapid thermal annealing (RTA) at 850° C. Immediately after the HCl dip to remove native oxide, samples were loaded in the atomic

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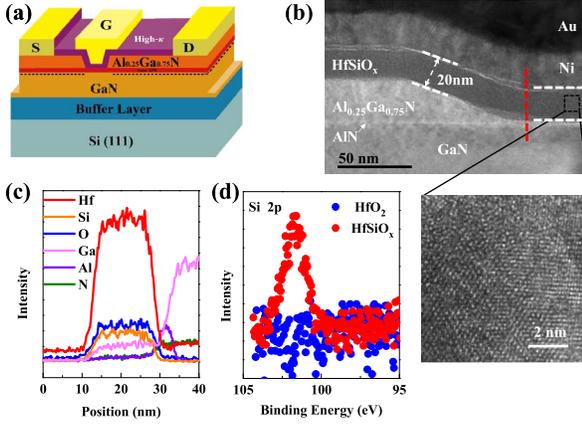


Fig. 1. (a) Cross section schematic of the GaN MOS-HEMT. (b) Cross-sectional TEM image of the recessed GaN MOS-HEMT in the recessed gate area. Inset is the HRTEM of the HfSiO_x dielectric. (c) The EDX analysis of heterointerfaces indicated by the red dashed lined in (b). (d) The XPS spectra in the Si 2p photoemission peak region of the HfO₂ and HfSiO_x surface.

layer deposition (ALD) chamber for gate dielectric deposition. A 20 nm HfO₂ dielectric was grown using precursors including Tetrakis(dimethylamino)hafnium and ozone at a deposition temperature of 300°C. As for the 20 nm HfSiO_x, one cycle reaction of Tris(dimethylamino)silane with ozone was inserted into HfO₂ by after 8 cycles of HfO₂ to form ~1 nm HfSiO_x. Devices were completed by Ni/Au gate metal deposition. Besides, the recessed HEMTs without oxide layer were also fabricated using the same process as a control sample. The HRTEM of HfSiO_x in the inset in Fig. 1(b) shows the dielectric is polycrystalline after the final device processing steps. The energy dispersive X-ray (EDX) analysis of the HfSiO_x oxides in Fig. 1(c) confirms Si atoms are uniformly distributed in HfSiO_x. Fig. 1(d) depicts the XPS spectra of Si 2p core level from 20 nm HfO₂ or HfSiO_x on recessed GaN epilayer. In the case of HfSiO_x, a peak can be discerned at ~102 eV, confirming the Si atoms in the HfSiO_x dielectric, while this peak is absent in HfO₂.

III. RESULTS AND DISCUSSION

Fig. 2(a) compares the transfer characteristics of HEMTs in saturation region. The recessed HEMT shows V_{tsat} of 0.75 V and SS of 133 mV/dec. The $g_{m,\text{max}}$ of HEMT is 109 mS/mm, which is lower than MOS-HEMT since it is overwhelmed by the gate leakage. The HfO₂MOS-HEMT exhibits increased V_{tsat} of 2.7 V, owing to the leakage blocking capability, but it shows large SS of 129 mV/dec. With the incorporation of silicon in HfO₂, the device performance has been significantly improved with near-ideal SS of 65 mV/dec, indicating lower interface trap density. The V_{tsat} shifts to 1.5 V, which can be attributed to less acceptor-like traps in HfSiO_x compared with HfO₂. Both MOS-HEMTs exhibit a high on/off ratio above 10^{10} . The HfSiO_x MOS-HEMT exhibits negligible hysteresis between forward and reverse sweep, which is smaller than HfO₂MOS-HEMT, indicating the suppression of interface traps.

To further understand the correlation between gate dielectrics and device performance, multi-frequency conductance method was applied to the FATFET structure with $L_g/W_g = 100/100 \mu\text{m}$ following the same device fabrication process as MOS-HEMTs [22], [23]. Fig. 2(b) shows the

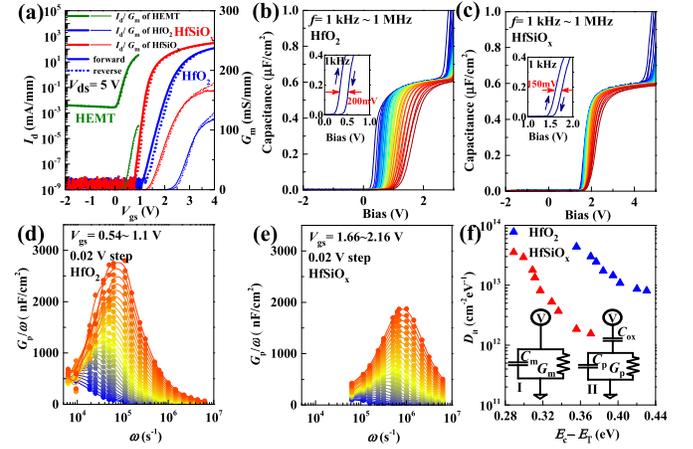


Fig. 2. (a) Double sweep transfer characteristics of HEMTs with 20 nm HfSiO_x, with 20 nm HfO₂ and without gate dielectric in saturation region at V_{ds} of 5 V. All HEMTs share the same dimensions: $L_g/L_{gs}/L_{gd} = 2/3.5/16 \mu\text{m}$. Double sweep CV characteristics of FATFETs (b) with 20 nm HfO₂ and (c) HfSiO_x. Frequency dependent conductance as a function of radial frequency for FATFETs (d) with 20 nm HfO₂ and (e) HfSiO_x. (f) The interface trap density distribution for FATFETs with HfO₂ and HfSiO_x. Inset model (I) is the equivalent circuit model of the FATFET, C_m and G_m are measured capacitance and parallel conductance. Model (II) is the simplified parallel model of the MOS model. The C_{ox} is the oxide capacitance and C_p and G_p are the equivalent parallel capacitance and conductance. All FATFETs share the same dimensions: $L_g/W_g = 100/100 \mu\text{m}$.

multi-frequency CV characteristics of HfO₂ FATFET with clockwise hysteresis ~ 200 mV at 1 kHz, which is slightly larger than 150 mV of HfSiO_x FATFET in Fig. 2(c). The frequency dependent V_{fb} for HfSiO_x FATFET is only 0.6 V, which is much smaller than 1.5 V for HfO₂ FATFET. At lower frequencies, the capacitance starts to increase from interface traps response, while at higher frequencies, slow traps cannot respond and the capacitance comes primarily from fast traps, which indicates large amount of slow traps located within HfO₂ or at the HfO₂/III-N interface, compared with HfSiO_x.

The gate-to-channel capacitance and conductance were measured with high terminal on the gate and low terminal on the source and drain of the FATFET, as shown in the model I in Fig. 2(f). The parallel capacitance C_p and conductance G_p can be extracted by subtracting the oxide capacitance, as shown in the model II. And the interface trap density can be extracted by the equation $D_{\text{it}} = 2.5(G_p/\omega)_{\text{peak}}/q$, where q is the electron charge [23], [24]. The trap energy level is given by $\Delta E = k_B T \ln(N_c \sigma_n v_t \tau_T)$, where k_B is the Boltzmann constant, T is the temperature, N_c is the effective conduction band DOS, σ_n is the capture cross-section, v_t is the thermal velocity of electrons, τ_T is the interface trap time constant corresponding with $2/\omega_c$ [23], [24]. Fig. 2(d) shows that G_p/ω peak of HfO₂/III-N interface gradually increases with ω , indicating higher interface trap density for interface traps with shorter emission constant. The G_p/ω peak movement of HfSiO_x/III-GaN interface in Fig. 2(e) reveals the same trend, and that the characterized interface traps by conductance method exhibit shorter emission time constant between 1 μs and 10 μs compared with HfO₂/III-N interface. As shown in Fig. 2(f), the minimal D_{it} extracted at $E_c - E_t = 0.37 \text{ eV}$ is $1.6 \times 10^{12} \text{ cm}^{-2} \text{ eV}^{-1}$ at HfSiO_x/III-N interface, much lower than the minimal D_{it} of $8.1 \times 10^{12} \text{ cm}^{-2} \text{ eV}^{-1}$ of HfO₂/III-GaN interface at $E_c - E_t = 0.43 \text{ eV}$. These D_{it} values extracted from conductance method are in accordance with SS obtained from

subthreshold region of I_D - V_{GS} curve, confirming better interface properties for HfSiO_x MOS-HEMT. It is supposed that the acceptor-like traps within the HfO_2 can be suppressed by SiO_2 interlayer [25]. Other works also show the incorporation of SiO_2 into high- κ dielectrics can screen the soft optical phonons associated with the highly polarizable bonds in high- κ dielectric HfO_2 [26].

In the multi-frequency CV measurement, only trap states with characteristic frequencies lying within 1 kHz \sim 1 MHz will contribute the single type interface trap signature during the measurement. Deep level traps within the oxide cannot follow the lowest gate signal frequency. In order to further characterize trap energy level within dielectrics, high- κ /III-N MIS structures were characterized by capacitance-mode deep-level transient spectroscopy. For both MISCAPs, the reverse bias was selected at pinch-off region, and the filling pulse bias was applied in strong inversion to fill the traps within the dielectrics. As shown in Fig. 3(a), one peak was found in HfO_2 MISCAP, indicating one electron trap level in HfO_2 . While two lower peaks were observed in HfSiO_x , which indicates two electron trap levels. The activation energy and capture cross-section can be obtained from Arrhenius plot based on the equation $\ln(e/IT^2) = \ln(16\pi\sigma_n k_B^2 m_n/h^3) - \Delta E_T/k_B T$, where e is the electron emission rate, m_n is the effective electron mass and h is the plank constant [27]. Fig. 3(b) shows the Arrhenius plot, the slope yields an energy level of 0.41 eV in HfO_2 , while there exist two shallower trap levels in HfSiO_x at 0.3 eV and 0.07 eV. The capture cross section of the trap in HfO_2 is $1.7 \times 10^{-11} \text{ cm}^2$, while the traps cross sections in HfSiO_x are 4.4×10^{-11} and $2.1 \times 10^{-17} \text{ cm}^2$. The trap concentration N_T can be obtained from $N_T = 2\Delta C/C_0 N_A$, where N_D is the net donor concentration obtained from CV profiling, ΔC is the difference between the capacitance when the traps are filled and emptied, and C_0 is the steady state capacitance [27]. The trap concentration of HfO_2 is $4 \times 10^{12} \text{ cm}^{-3}$, which is almost twice the number in HfSiO_x with trap concentrations of $2.7 \times 10^{12} \text{ cm}^{-3}$ and $2.1 \times 10^{12} \text{ cm}^{-3}$. Based on the CV analysis and DLTS analysis, the defects properties in HfO_2 and HfSiO_x are uncovered and compared. As shown in Fig. 3(c), the highly polarizable bonds in high- κ dielectric HfO_2 induces large number of oxide traps as well as interface traps. With the incorporation of SiO_2 into high- κ dielectrics, two shallower trap levels are formed, as shown in Fig. 3(d). The new trap levels can screen the soft optical phonons associated with the polarizable bonds in HfO_2 , and mitigate its interaction on dielectric/III-N interface and channel mobility [26].

To evaluate the reliability of these MOS-HEMTs, three groups of devices were stressed at constant gate-source voltage with source and drain grounded for each MOS-HEMT. The time when the current suddenly increases is defined as time-to-breakdown t_{BD} . Fig. 4(a) shows the time-to-breakdown is Weibull distributed for HfO_2 MOS-HEMTs and linear fitting was used for the full range. The Weibull slope of HfO_2 decreases from 4.2 to 1.5 with increasing oxide field. The oxide breakdown can be ascribed to field-induced defects. Under high oxide field, the traps needed to form the conductive breakdown path are randomly redistributed within the oxide and are broad spread statistically, which results in smaller Weibull slope [28]. Fig. 4(b) shows the Weibull slope of HfSiO_x MOS-HEMTs decreases from 3.8 to 2.0 with increased oxide field. The analysis of time-to-breakdown at 63.2% failure

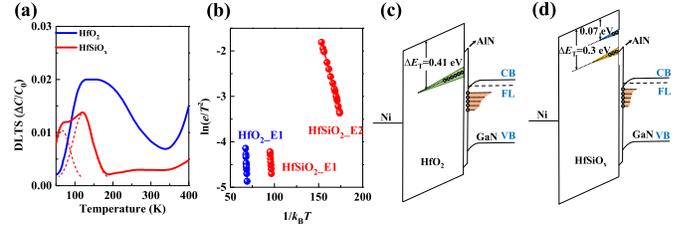


Fig. 3. (a) C-DLTS spectra of recessed GaN MISCAPs with 20 nm HfSiO_x and 20 nm HfO_2 at $t_1/t_2 = 2 \text{ ms}/15 \text{ ms}$. (b) The Arrhenius plot for the deep level traps for both MISCAPs. One electron trap was observed in HfO_2 , while two electron traps were observed in HfSiO_x . Band diagram of GaN MOS-HEMT devices with (c) 20 nm HfO_2 and (d) 20 nm HfSiO_x at positive gate-source bias.

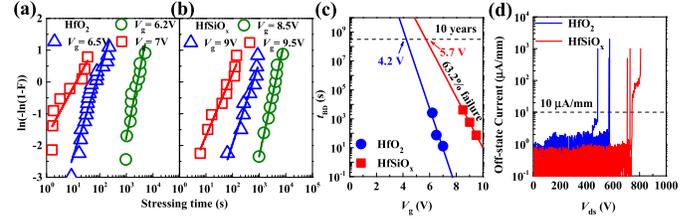


Fig. 4. Weibull distribution of time-to-breakdown for (a) HfO_2 and (b) HfSiO_x MOS-HEMTs. The Weibull slopes of HfO_2 MOS-HEMTs are 4.2, 2.7, and 1.5 for $V_g = 6.2 \text{ V}$, 6.5 V and 7 V, respectively. The Weibull slopes of HfSiO_x MOS-HEMTs are 3.8, 2.5, and 2.0 for $V_g = 8.5 \text{ V}$, 9V and 9.5V, respectively. (c) Lifetime prediction of 63.2% failure level for HfO_2 and HfSiO_x gate dielectrics. (d) Off-state characteristics of several HfO_2 and HfSiO_x MOS-HEMTs with the same dimensions at gate-source bias of 0 V. The dimensions are $L_g/L_{gs}/L_{gd} = 2/3.5/16 \text{ }\mu\text{m}$.

rate (i.e. $\ln(-\ln(1-F)) = 0$) for each dielectric allows one to predict the lifetime, as shown in Fig. 4(c) [29]. The maximum gate bias for ten-year lifetime is extracted to be 5.7 V ($\sim 2.85 \text{ MV/cm}$) for HfSiO_x MOS-HEMT and 4.2 V ($\sim 2.1 \text{ MV/cm}$) for HfO_2 MOS-HEMT at 63% failure level. The better time-to-breakdown of the HfSiO_x MOS-HEMT can be attributed to the better dielectric interface as well as the larger energy band offset of HfSiO_x on III-N compared with HfO_2 which effectively suppress gate leakage current and increase the electric field strength [30]. The reversible breakdown property in ultra-thin dielectrics is not observed in power MOS-HEMTs with thick dielectrics [31].

Fig. 4(d) shows the off-state breakdown voltage for both E-mode MOS-HEMTs. The reduced interface defects and high electric strength of HfSiO_x can enhance the off-state breakdown voltage as well. For HfO_2 MOS-HEMTs, drain-source current starts to increase after a drain-source bias of 400 V, then sharply increase to $10 \text{ }\mu\text{A/mm}$ with a maximum drain-source bias $\sim 580 \text{ V}$. With suppressed gate leakage current, the off-state current of HfSiO_x MOS-HEMT remains below $1 \text{ }\mu\text{A/mm}$ up to a drain-source voltage of 700 V, resulting in a much higher off-state breakdown voltage of 742 V.

IV. CONCLUSION

In conclusion, this letter reports high- κ dielectrics in recessed gate E-mode GaN MOS-HEMTs, and silicon inclusion into HfO_2 dielectric provides better passivation in the recess channel region and minimized interface and fixed oxide traps and near-ideal subthreshold slope can be achieved. The dielectric breakdown strength of the E-mode GaN MOS-HEMTs can also be improved using HfSiO_x as gate dielectric compared to that based on HfO_2 .

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